



# Double O–Ne–Mg white dwarfs merging as the source of the powerful gravitational waves for LIGO/VIRGO type interferometers



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## HIGHLIGHTS

- There are new strong non spiralling gravitational waves source for LIGO/VIRGO detectors.
- It is noted that double O–Ne–Mg white dwarfs mergers can produce strong gravitational waves with frequencies in the 1 kHz range.
- Such events can be followed by the SuperNova type Ia phenomena and can be registered by LIGO/VIRGO detectors.

## ARTICLE INFO

### Article history:

Received 28 February 2017

Revised 25 April 2017

Accepted 26 April 2017

Available online 27 April 2017

### Keywords:

Gravitational waves  
Gravitational collapse  
Double white dwarfs  
GW  
LIGO/VIRGO

## ABSTRACT

New strong non-spiralling-in gravitational wave (GW) source for LIGO/VIRGO detectors are proposed. Double O–Ne–Mg white dwarf mergers can produce strong gravitational waves with frequencies in the several hundreds Hz range. Such events can be followed by a Super Nova type Ia.

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The discovery of gravitational waves (Abbot et al., 2016a,b,c) produced by a black-hole merger and predicted by the first Monte-Carlo code (Kornilov and Lipunov, 1983; Lipunov, 1992; Lipunov et al., 1996b) of the ScenarioMashine (Lipunov et al., 1997a,b,c) confirmed the most general ideas about the evolution of baryonic matter in the Universe in the form of the evolution of binary stars (van den Heuvel and Heise, 1972; van den Heuvel and Loore, 1973; Tutukov and Yungelson, 1973, Tutukov and Masevich, 1980; Tutukov and Yungelson, 1993; Lipunov et al., 2017). The next task is to discover mergers involving two neutron stars or a neutron star and a black hole, see references Clark et al. (1979), Lipunov (2016). However, we will show that there is yet another channel for the generation of powerful gravitational waves during the formation of a massive neutron star or a light black hole and this event should be accompanied by a SNIa explosion.

Because of the idea suggested by Iben and Tutukov (1984) and Webbink (1984) we now know that a merger of white dwarfs may result in an SNIa event. Moreover, SNIa observations in elliptical galaxies (Totani et al. (2008) demonstrate that white-dwarf mergers are the main channel for the formation of SNIa in elliptical galaxies, at least after the first billion years of their life time (Lipunov, Panchenko, and Pruzhinskaya, 2011).

The condition for explosion is that the total mass of white dwarfs must exceed the Chandrasekhar limit, see Masevich and Tutukov (1981); van den Heuvel (2011); Livio (2000):

$$M_{\text{Ch}} \sim 5.83 \mu_e^{-2} \sim 1.4 M_{\odot}, \quad (1)$$

where  $\mu_e \approx 2$  is the average nucleon to electron ratio in the white dwarf.

White-dwarf mergers must be accompanied by a gravitational-wave pulse (Cutler and Thorne, 2002). It was, however, pointed out that the low amplitude and low frequency ( $\sim 1$  Hz) of such a pulse compared to those of pulses accompanying NS and BH mergers these events are of little interest for LIGO type interferometers, at least in the years to come.

The point is that the overwhelming majority of merging white dwarfs whose merger results in a SNIa explosion are the so called CO-type white dwarfs. After the merger an unstable CO-type dwarf forms, which, after the beginning of the collapse, ignites, breaks apart, and scatters as a result of a thermonuclear explosion leaving no remnant, see Iben and Tutukov (1984), van den Heuvel (2011). The maximum amplitude of the GW at the merger point can actually be estimated as  $h_0 \sim R_g/R_{\text{WD}} \sim 10^{-3}$ , which is about 100 times less than expected in the case of a double neutron star

merger, resulting in a factor of  $10^6$  times smaller detection volume. Furthermore, the characteristic rotation frequency of a white dwarf should be, as we already wrote, well below the sensitivity curve of interferometers.

However, a merger of the more massive O–Ne–Mg dwarfs results in quite a different picture. After the merger of the dwarfs an O–Ne–Mg core may form with a mass exceeding 2.5 solar masses. Such a white dwarf will have huge angular momentum and the collapse may result in the development of a Dedekind type instability and formation of large quadrupole moment. Unlike the collapse of merged CO dwarfs that of merged O–Ne–Mg dwarfs does not result in complete disintegration (Nomoto, 1984; Tominaga et al., 2013), and given such a large mass it would be natural for a massive neutron star or, more likely, a black hole to form. In this case the initial GW amplitude should be about  $h_0 \sim 0.1$ – $0.3$ , which is comparable to the case of a usual merger of two relativistic stars.

The collapse of O–Ne–Mg dwarfs may also proceed via a different way, in accordance with the so called SD mechanism of supernova formation. An O–Ne–Mg dwarf may build up mass as a result of accretion in a binary system (usually with a red dwarf companion).

This is the so called Accretion Induced Collapse. However, this scenario results in a much smaller angular momentum of the final neutron star or black hole. An extreme white dwarf with a higher than critical mass is a very centrally concentrated object with a density contrast of  $\sim 1/50$ . Hence most of the angular momentum should be distributed in the exterior layers of the white dwarf whereas its more massive central part would have much smaller angular momentum and hence possesses a small quadrupole moment.

Thus a merger of two O–Ne–Mg dwarfs provides a competing channel for the generation of strong bursts of gravitational waves compared to NS and BH mergers. The evolutionary sequences:

$$WD_{O-Ne-Mg} + WD_{O-Ne-Mg} > NS + SNIa + GWB$$

or

$$WD_{O-Ne-Mg} + WD_{O-Ne-Mg} > BH + SNIa + GWB$$

depend on the Openheimer-Volkoff limit value. The collapse of a heavy no-solid-body-rotating O–Ne–Mg WD is a very complex process.

The maximal frequency is equal to 2 times the maximal observed spin frequency for a NS. The maximal spin frequency is well known from millisecond pulsar observations and is equal about 700 Hz. This maximal frequency may be connected with the neutron star equation of state (Lipunov and Postnov, 1984; Friedman et al., 1985). In the case of NS formation, a fast rotating pulsar can be observed.

Let us estimate the probability of these events. The mass ratio distribution of the main-sequence progenitors is  $\varphi(\mathbf{q}) = \text{const}$  (Lipunov et al., 1996a, Fedorova et al., 2004), where  $0 < \mathbf{q} = M_2/M_1 < 1$ .

In our case both components should have masses in the  $M_0 \sim 8$ – $10 M_\odot$  interval. The probability of the formation of a binary of two stars of similar mass within 10% is about  $\sim 10$  times lower than the general probability of the formation of a binary stars. However, we need two sufficiently massive stars. Such stars form in accordance with the Salpeter function:

$$dN/dM \sim M^{-\alpha} \quad (2)$$

Where  $\alpha = 2.35$ . I assume all stars more massive than 3 solar masses up to 8 solar masses to form CO white dwarfs sufficiently massive to produce a Type Ia supernova, and all stars between 8 and 10 solar masses to form Type Ia supernovae due to the merger of O–Ne–Mg white dwarfs.

The relative probability of the formation of two O–Ne–Mg dwarfs compared to that of the formation of binary white dwarfs with total mass more than the Chandrasekhar limit is of about  $10^{-1}$ . We can therefore expect  $10^{-2}$  of all type Ia supernovae possibly to be strong gravitational wave sources.

The SNIa frequency at the current epoch, i.e. after about 10 Gyr (see Jørgensen et al., 1997; Totani, 2008; Lipunov et al., 2011), is  $\sim 1/300 \text{ yr}^{-1}$  per  $10^{11}$  solar masses. At the same time, the frequency of double NS mergers is  $\sim 1/10,000$ – $1/30,000$  per  $10^{11}$  solar mass (Lipunov et al., 1987; Grishchuk et al., 2001; Kalogera et al., 2004). Hence the collapse rate of binary O–Ne–Mg dwarfs may occur with about the same frequency as double neutron star mergers, and given that we are dealing with collapse into a black hole whose GW amplitude is twice higher than for neutron stars the corresponding detection rate may be even higher than the detection rate of double neutron star mergers.

I am grateful Alexander Tutukov and Edward van den Huevel for usefull remarks. This work was supported in part by the Development Program of Lomonosov Moscow State University, Moscow Union OPTICA, Russian Science Foundation 16-12-00,085.

## References

- Abbott, B., et al., 2016a. *ApJ* 818L, 22A.  
 Abbot, B., et al., 2016b. *Phys. Rev. Lett.* 116, 1102.  
 Abbot, B., et al., 2016c. *ApJL* 826, 13A.  
 Clark, J.P.A., van den Huevel, E.P.J., Sutantyo, W., 1979. *Astron. Astrophys.*  
 Cutler, C., Thorne, K., 2002. *grg.conf.*, p. 72C.  
 Friedman, J.L., Ipser, J.R., Parker, L., 1985. *ApJ* 292, 111–117.  
 Fedorova, A.V., et al., 2004. *Astron. Lett.* 30, 73.  
 Grishchuk, L.P., Lipunov, V.M., Konstantin, P., Prokhorov, A., Mikhail, E., Sathyaprakash, B.S., 2001. *Phys. Uspekhi* 44, 1.  
 van den Heuvel, E.P.J., Heise, J., 1972. *Nature. Phys. Sci.* 239, 67.  
 van den Heuvel, E.P.J., Loore, D., 1973. *Astron. Astrophys* 25, 387–395.  
 van den Heuvel, E.P.J., 2011. *BASI* 39, 1V.  
 Jørgensen, H.E., Lipunov, V.M., Panchenko, I.E., et al., 1997. *ApJ* V 486 (1), 110–116.  
 Iben, I., Tutukov, A.V., 1984. *ApJS* 54, 335.  
 Kalogera, et al., 2004. *ApJ* 614, L137.  
 Kormilov, V., Lipunov, V., 1983. *SvA* 27, 334K.  
 Lipunov, V.M., 2016. *PhyU* 59, 918.  
 Lipunov, V.M., Postnov, K.A., Prokhorov, M.E., 1987. *A&A* 176L, 1L.  
 Lipunov, V., et al., 2017. *New Astronomy* 51, 122–127.  
 Lipunov, V., Panchenko, I., Pruzhinskaya, M., 2011. *New Astronomy* 16, 250L.  
 Lipunov, V.M., Panchenko, I.E., Pruzhinskaya, M.V., 2011. *New Astronomy* 16, 250–252.  
 Lipunov, V., Postnov, K., Prokhorov, M., 1996. *Astron. Astrophys.* 310, 489–507.  
 Lipunov, V.M., Postnov, K.A., 1984. *Astroph. Space. Sci.* 106, 103–115.  
 Lipunov, V., Postnov, K., Prokhorov, M., 1997a. *Astron. Lett.* 23 (4), 492–497.  
 Lipunov, V., Postnov, K., Prokhorov, M., 1997b. *New Astron.* 2 (1), 43–52.  
 Lipunov, V.M., Postnov, K., Prohorov, M.E., 1997c. *MNRAS* 288 (1), 245–259.  
 Lipunov, V., Postnov, K., Prokhorov, M., 1996. In: Sunyaev, R.A. (Ed.), *The Scenario Machine: Binary Star Population Synthesis*. Harwood Academic Publishers, Amsterdam.  
 Lipunov, V.M., 1992. *Astrophysics of Neutron Stars*. Springer, NewYork/Berlin.  
 Livio, M., 2000. In: Niemeyer, J.C., Truran, J.W., (Eds.), Cambridge University Press, Cambridge, 33.  
 Masevich, A.G., Tutukov, A.V., 1981. *INTSA*. 17.  
 Nomoto, K., 1984. *ApJ* 277, 791.  
 Tutukov, A.V., Masevich A., G., 1980. *Sov. Phys. Usp.* 23, 706.  
 Tominaga, N., Blinnikov, S., Nomoto, K., 2013. *ApJ*. 771L, 12T.  
 Totani, T., Morokuma, T., Oda, T., et al., 2008. *PASJ* 60, 1327.  
 Tutukov, A., Yungelson, L., 1973. *Nauchnye Informatsii* 27, 70.  
 Tutukov, A.V., Yungelson, L.R., 1993. *MNRAS* 260, 675T.  
 Webbink, R., 1984. *ApJ* 277, 355.